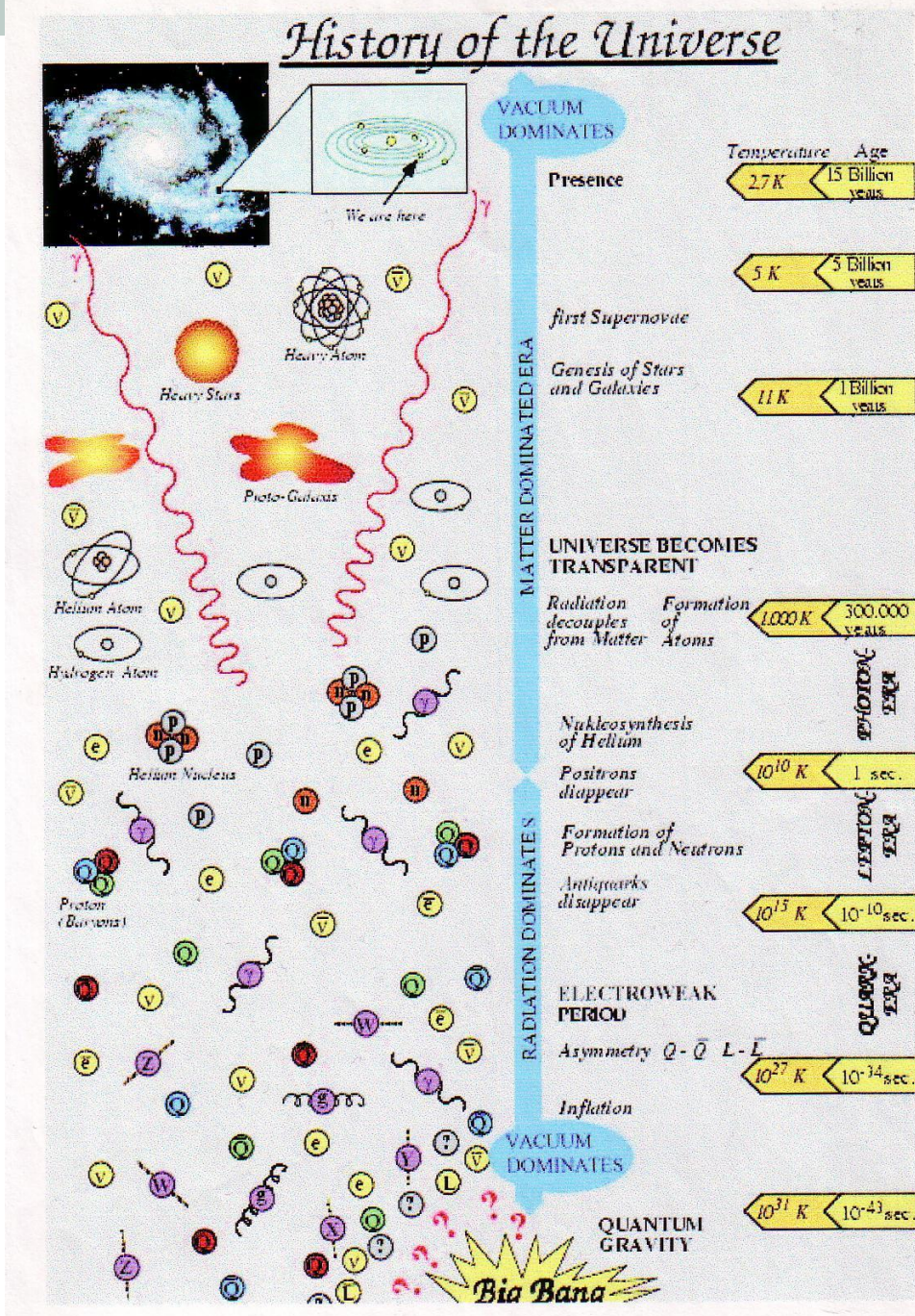


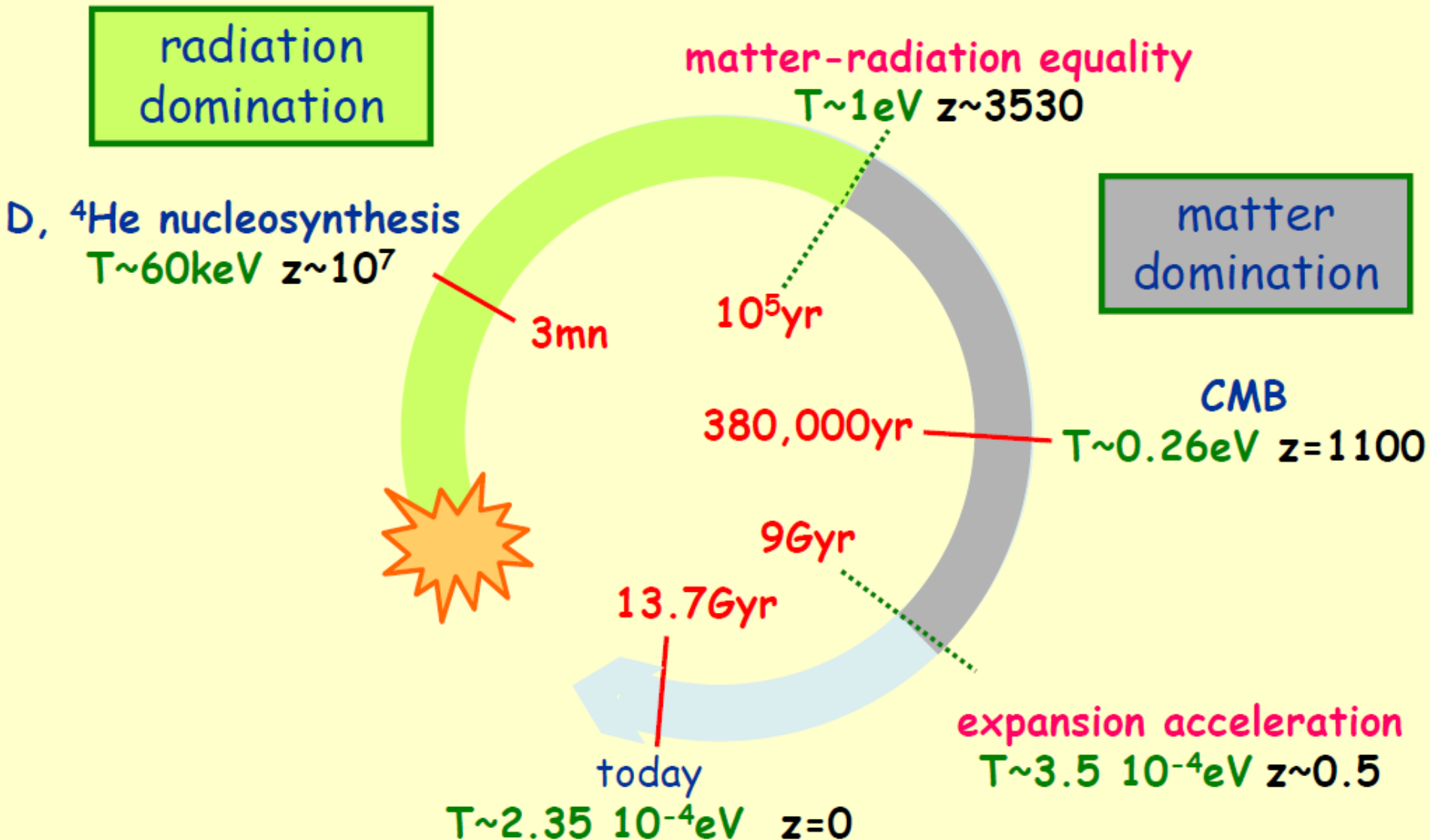
FIZIKALNA

KOZMOLOGIJA

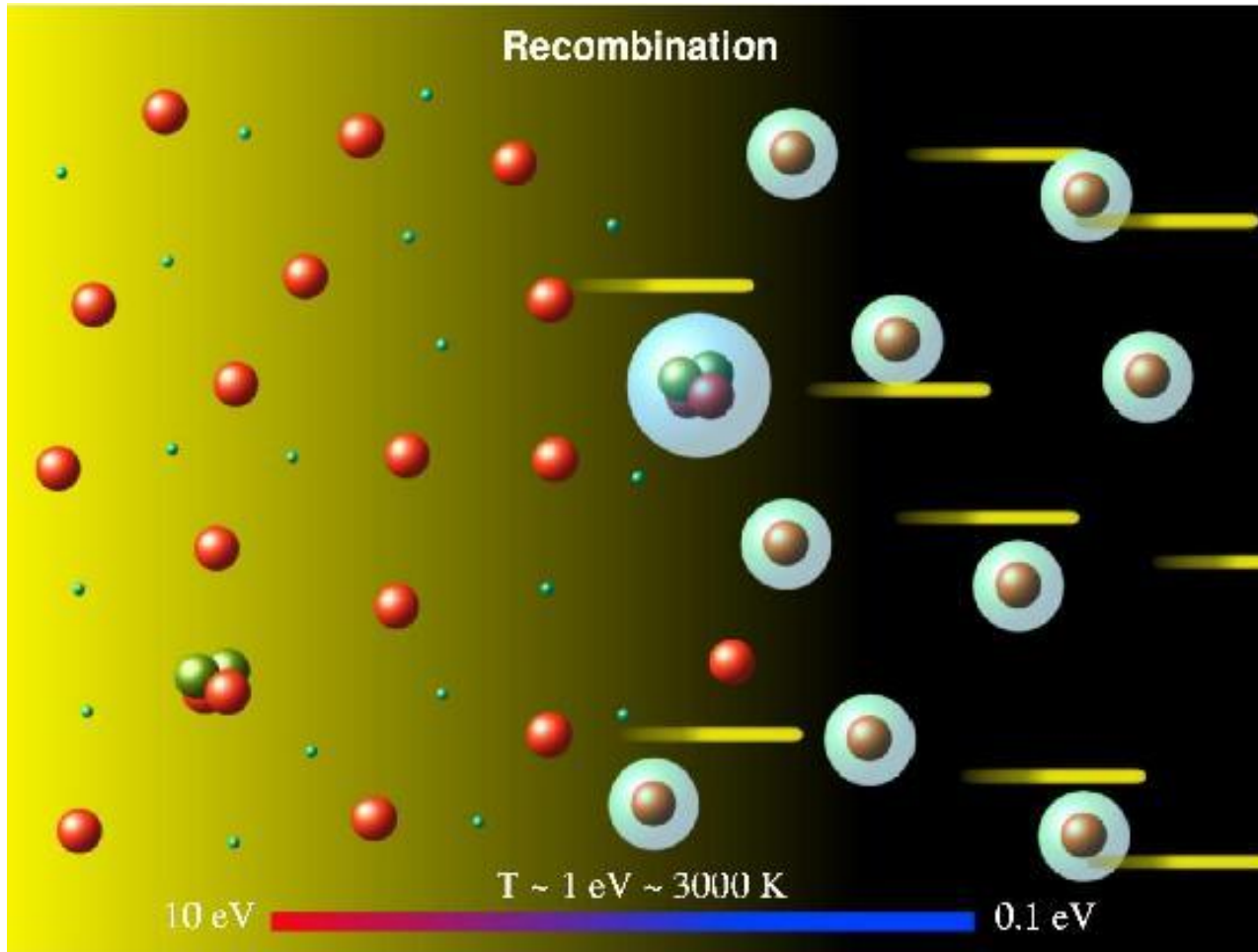
VIII. FIZIKA CMB-a

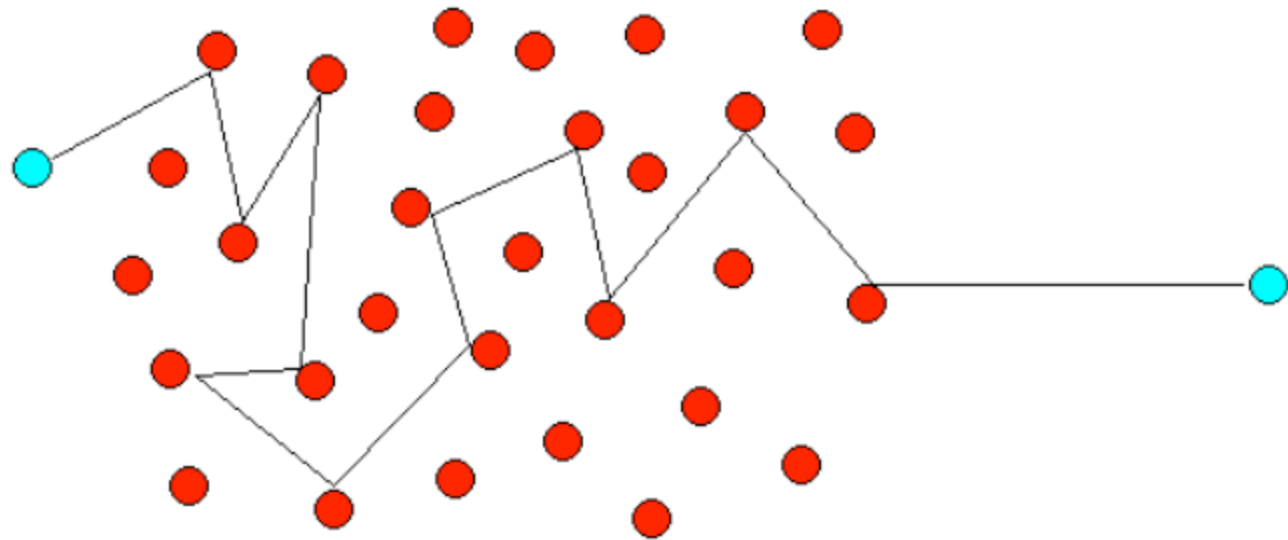


"KOZMIČKI SAT"



SVEMIR U ODREĐENOM TRENUTKU IZLAZI IZ FAZE PLAZME





the Universe is opaque at high densities (the mean free path of a photon is very short), as the density drops with time, the Universe becomes transparent (the mean free path of a photon becomes very large).

The ionisation fraction x_e drops rapidly at (re)combination so the Thomson scattering rate also decreases sharply below the Hubble expansion rate ... this defines a **last scattering surface** for the relic photons ... which we see today as the cosmic microwave background

Thomson scattering on electrons: $\gamma + e \rightarrow \gamma + e$

Interaction rate for photons:

$$\Gamma_{\text{Thomson}} = n_e \langle \sigma_T |v| \rangle \propto x_e T^3 \sigma_T$$

Expansion rate of the universe (MD era): $H \propto T^{3/2}$

$\Gamma_{\text{Thomson}} > H$: Photons and matter *in equilibrium*

$\Gamma_{\text{Thomson}} < H$: Photons and matter stop interacting

Na $T \sim 1\text{eV}$ fotoni vezani za e^- Compt.raspršenjem

Visok omjer fotona prema barionima $(\eta_b)^{-1} \rightarrow$ neutralnog H vrlo malo

Jedn.ravnoteže za $e^- + p \leftrightarrow H + \gamma$, uz $n_\gamma = n_\gamma^{(0)}$:

$$\frac{n_e n_p}{n_H} = \frac{n_e^{(0)} n_p^{(0)}}{n_H^{(0)}}$$

Neutralnost svemira: $n_e = n_p$

def.omjer slobodnih e^- :

$$X_e \equiv \frac{n_e}{n_e + n_H} = \frac{n_p}{n_p + n_H}$$

Uz $n_i^{(0)} = g_i \left(\frac{m_i T}{2\pi}\right)^{3/2} e^{-m_i/T}$ imamo:

$$\frac{X_e^2}{1 - X_e} = \frac{1}{n_e + n_H} \left[\left(\frac{m_e T}{2\pi}\right)^{3/2} e^{-[m_e + m_p - m_H]/T} \right] \quad \text{Saha}$$

energija ionizacije: $m_e + m_p - m_H = \epsilon_0 = 13.6 \text{ eV}$

$X_e \sim 1$ (sav H ioniziran) sve dok T ne padne na $T \ll \epsilon_0 \rightarrow$ rekobinacija
 \rightarrow Boltzm.jedn. daje evoluciju X_e

- Boltzmannova jedn. uz uvrštavanje $n_e^{(0)} n_p^{(0)} / n_H^{(0)}$ iz Sahine :

$$a^{-3} \frac{d(n_e a^3)}{dt} = n_e^{(0)} n_p^{(0)} \langle \sigma v \rangle \left\{ \frac{n_H}{n_H^{(0)}} - \frac{n_e^2}{n_e^{(0)} n_p^{(0)}} \right\} = n_b \langle \sigma v \rangle \left\{ (1 - X_e) \left(\frac{m_e T}{2\pi} \right)^{3/2} e^{-\epsilon_0/T} - X_e^2 n_b \right\}$$

Brzina ionizacije: $\beta \equiv \langle \sigma v \rangle \left(\frac{m_e T}{2\pi} \right)^{3/2} e^{-\epsilon_0/T}$, $\alpha^{(2)} \equiv \langle \sigma v \rangle$

Samo rekombinacija pobuđenog stanja, $\alpha^{(2)}$, daje doprinos

→ Boltzmannova jedn.:

$$\frac{dX_e}{dt} = \left\{ (1 - X_e)\beta - X_e^2 n_b \alpha^{(2)} \right\}$$

numeričko rješavanje

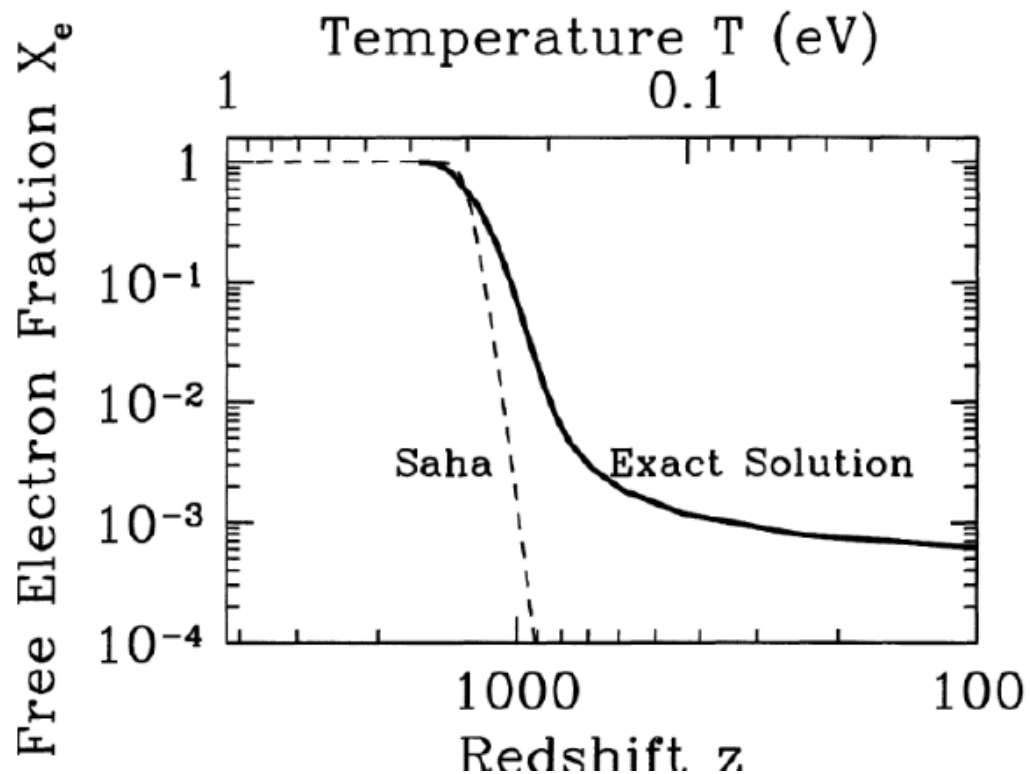
Rekomb. postaje značajna

na $T \sim \frac{1}{4}$ eV

to odgovara crvenom pomaku

$z \sim 1000$

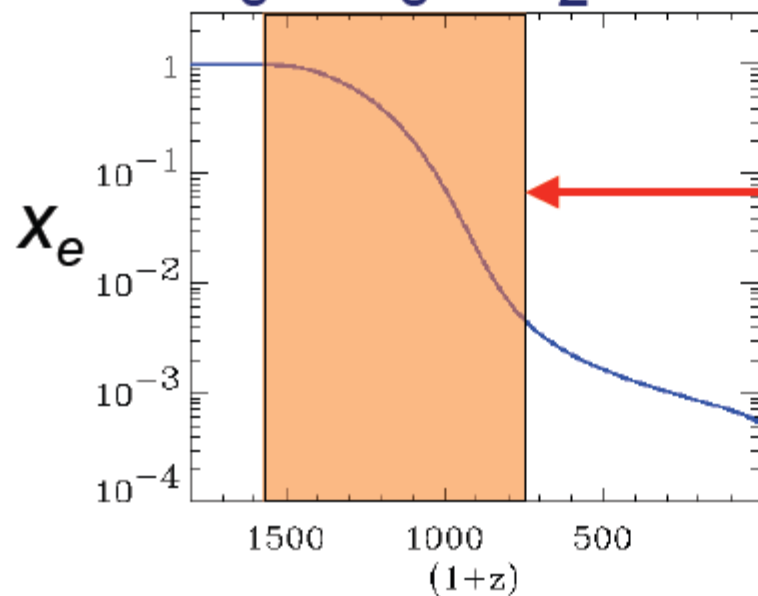
odvezivanje fotona



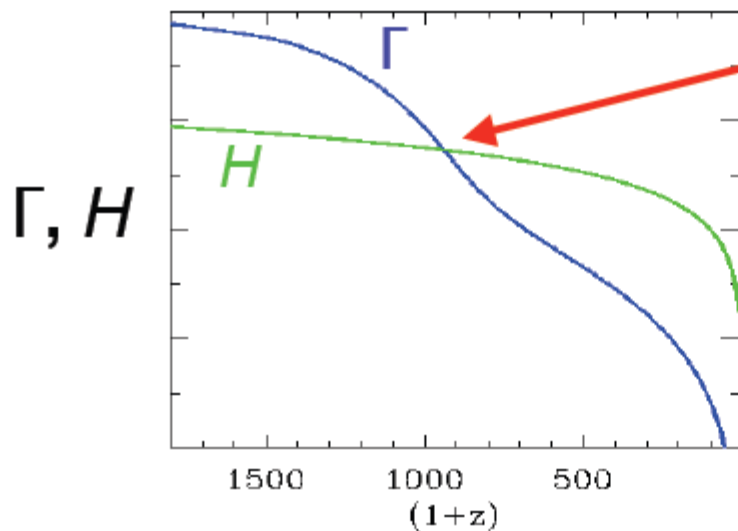
T/eV ←

0.3 0.2 0.1

5 3 2

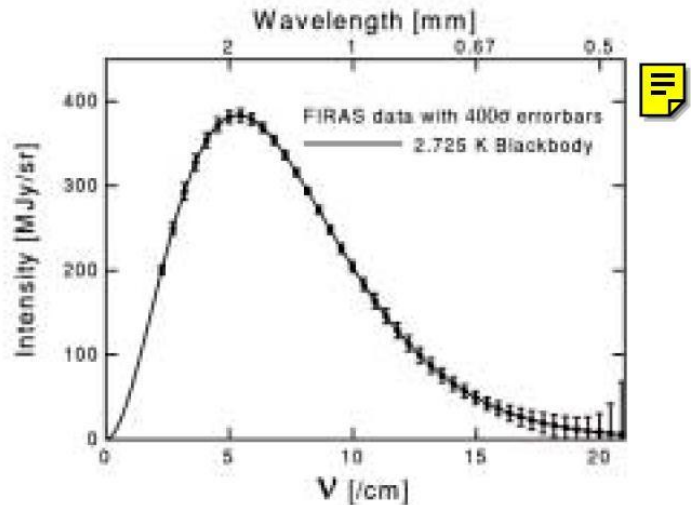


Recombination
(according to the
Saha ionisation eq.)

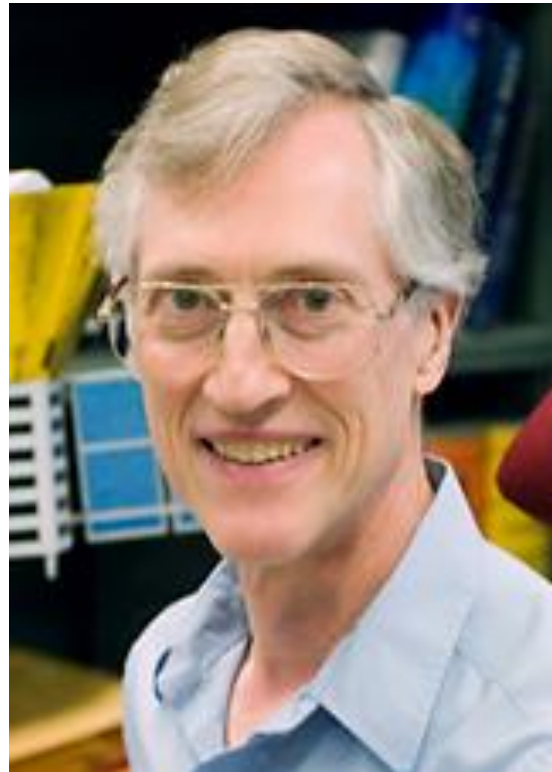
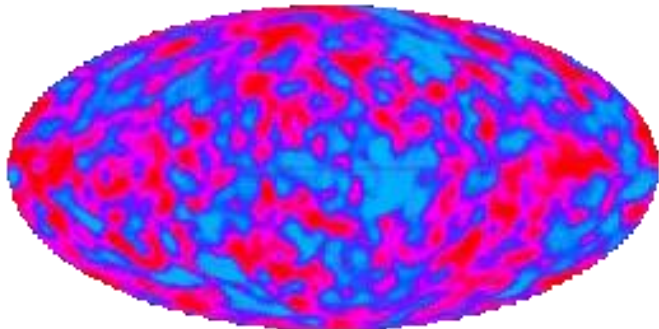


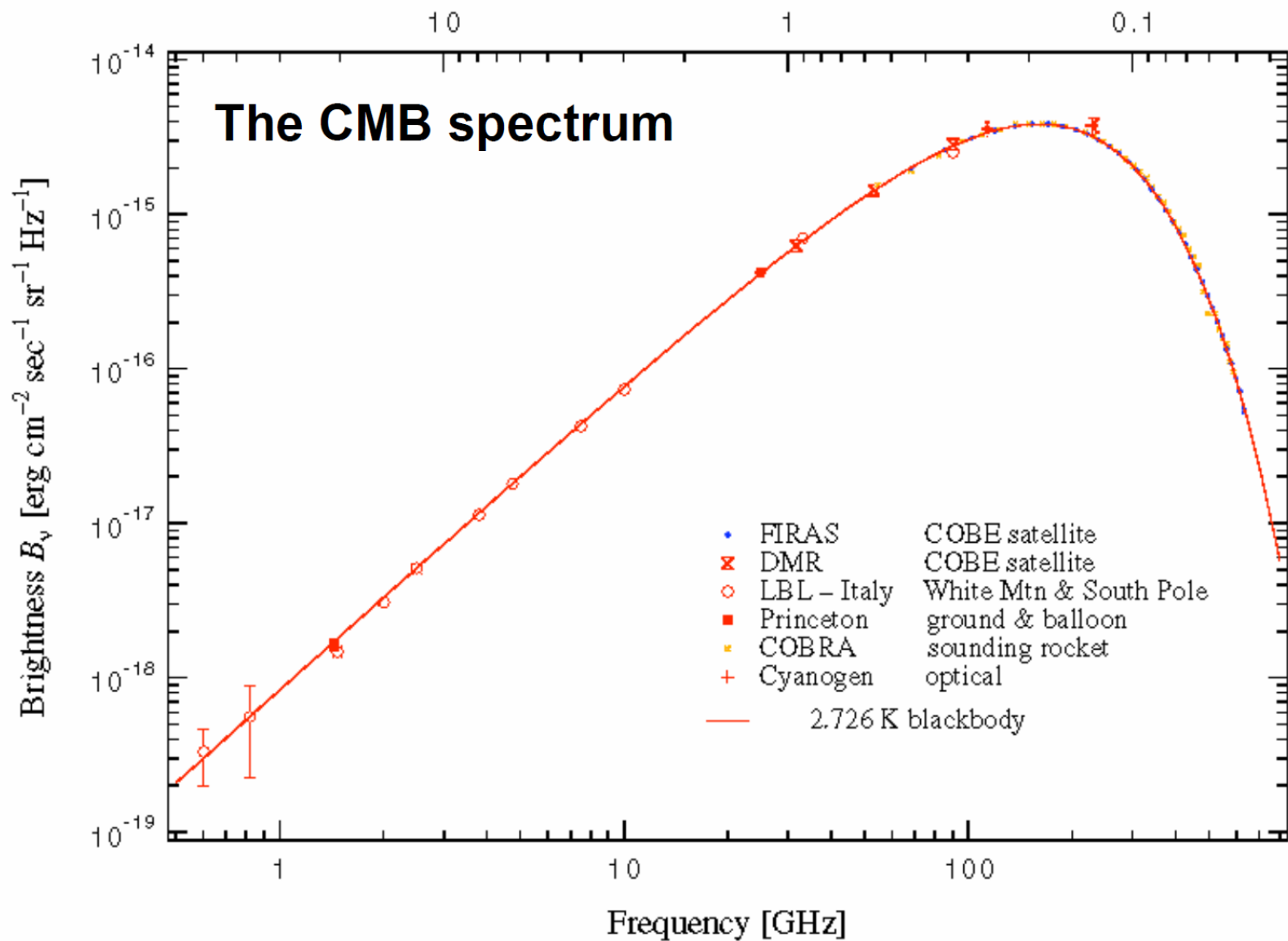
Decoupling of
photons and baryons

Nobelova nagrada iz fizike za 2006: "za otkriće prirode CMB-a" - John C. Mather & George F. Smoot (vidjeti MFL 2/226, str. 104)



FIRAS spectrum (frequency vs. intensity) with errors expanded 400 times





Such a perfect blackbody is testimony to our hot, dense past

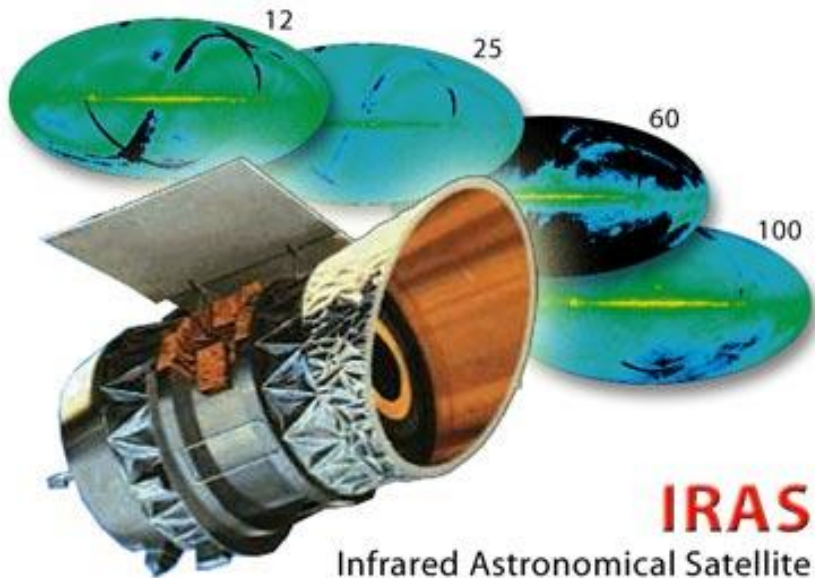
POZADINSKO ZRAČENJE (CMB)

- spektar zračenja crnog tijela,
za DZ:

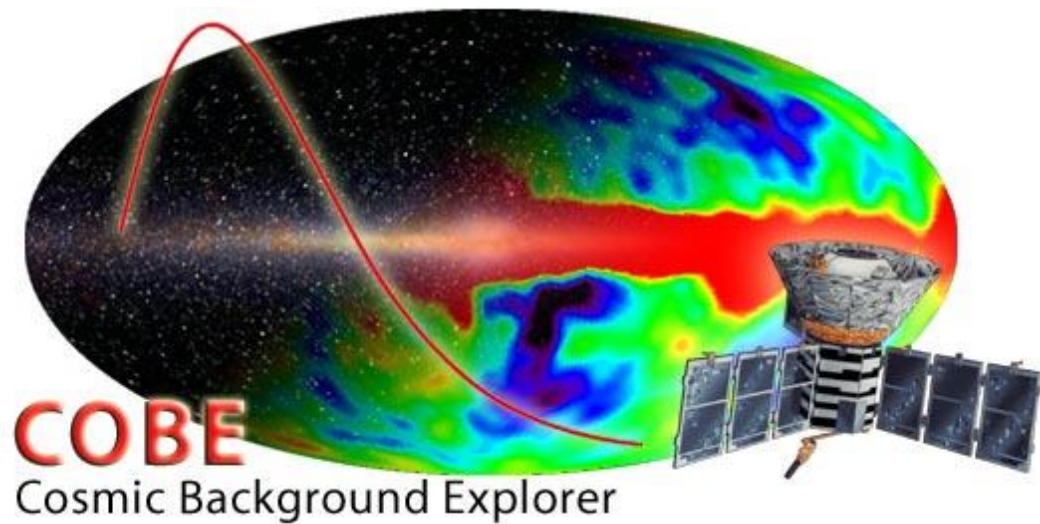
- **OBLIK SPEKTRA OČUVAN PRI EKSPANZIJI SVEMIRA**
- **SPEKTAR CRNOG TIJELA U SVAKOM INERCIJALNOM SUSTAVU**
- **PROBLEM HORIZONTA**

SATELITSKI POKUSI (MJERENJA POZADINSKOG ZRAČENJA)

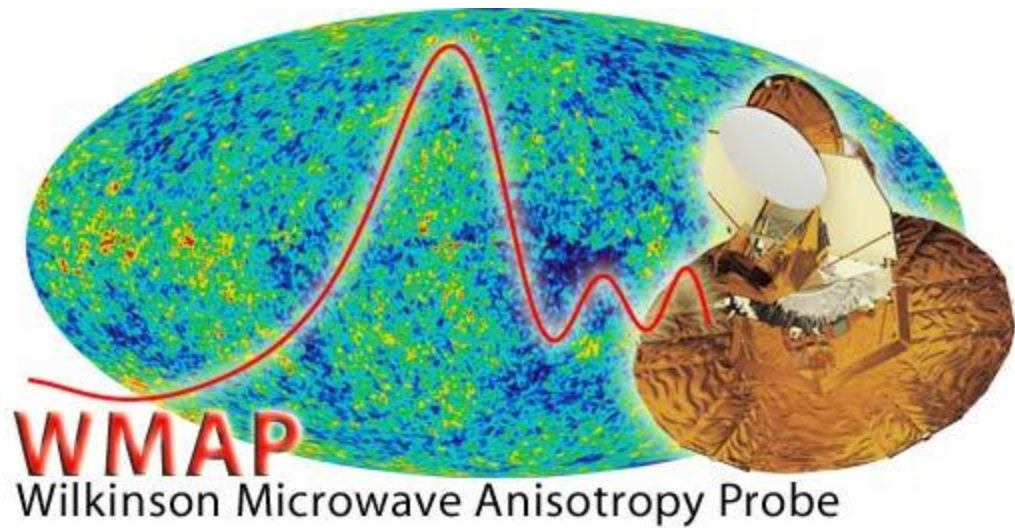
1983.



1989 – 93.



2001 –



POZADINSKO MIKROVALNO ZRAČENJE

1965



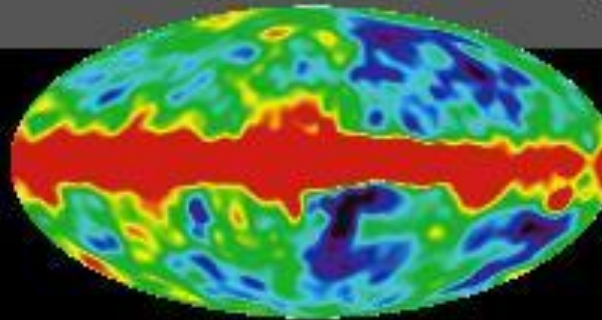
Penzias and
Wilson



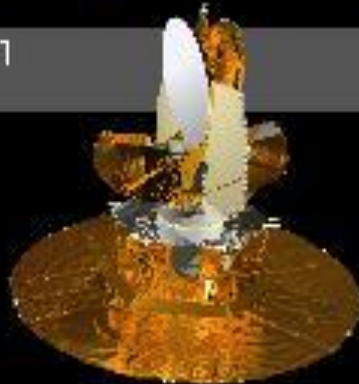
1992



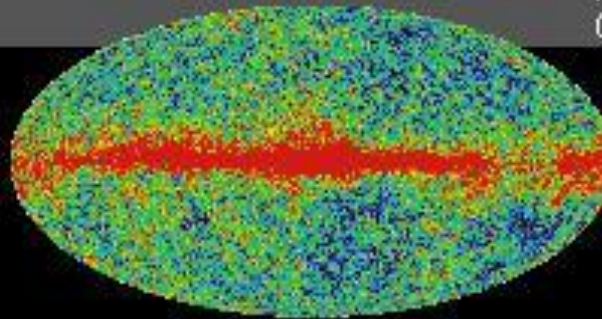
COBE



2001



MAP
(Simulated)



ESA lansirala Herschel Space Laboratory i PLANCK Satelite



DIPOLNA ANIZOTROPIJA CMB-a

Motion of the Earth wrt to the CMB retains exactly the thermal spectrum, but the temperature is boosted:

$$(1) \quad T_{\text{obs}} = \mathcal{D}T_0$$

where \mathcal{D} is the general form of the Doppler factor:

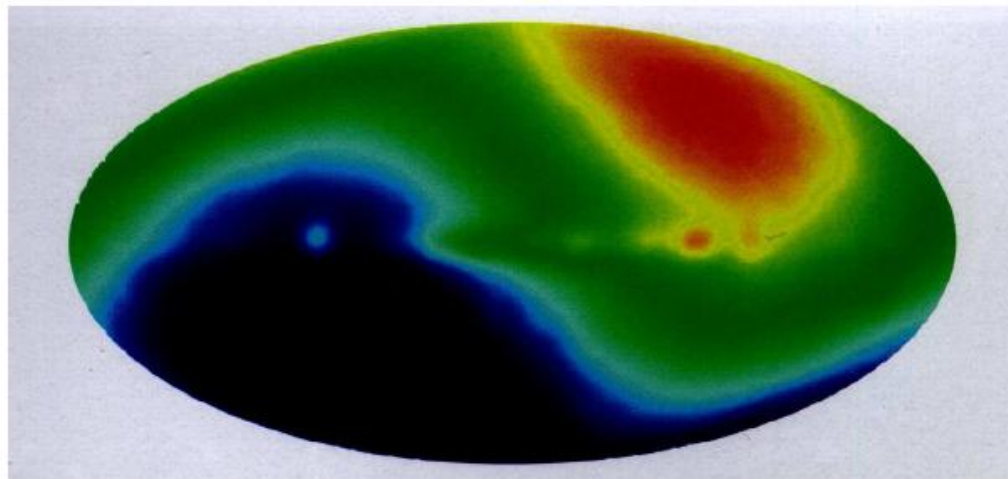
$$(2) \quad \mathcal{D} = \gamma^{-1} [1 - (v/c) \cos \theta]^{-1}$$

$$(3) \quad T_{\text{obs}} = \frac{T_0}{\gamma [1 - (v/c) \cos \theta]}$$

which can be expanded as a series of multipoles:

$$(4) \quad T_{\text{obs}} = T_0 \left[1 + \frac{v}{c} \cos \theta + \frac{1}{2} \left(\frac{v}{c} \right)^2 \cos 2\theta + O(v^3) \right]$$

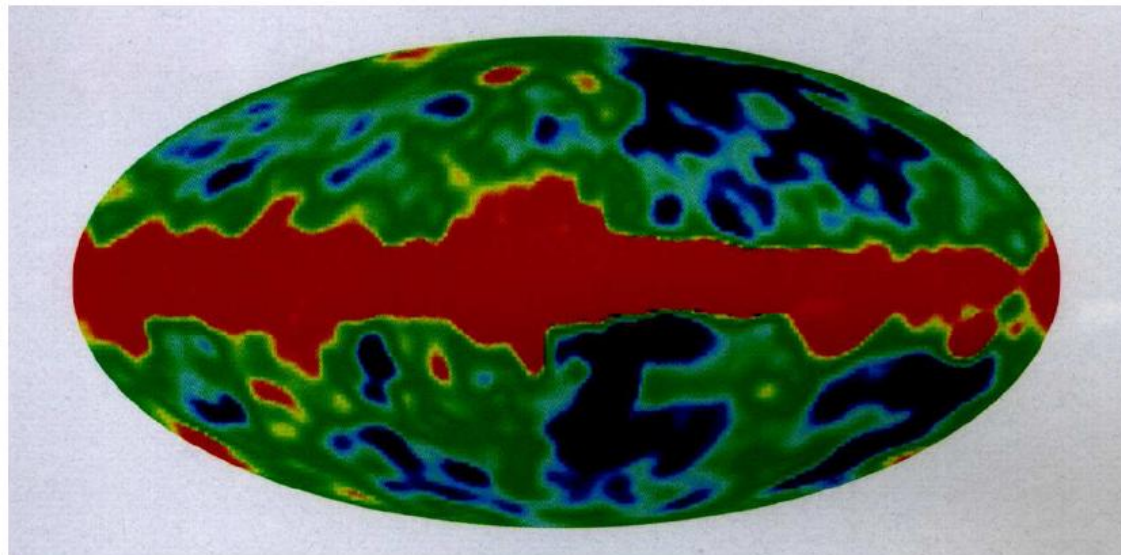
DIPOLNA ANIZOTROPIJA CMB-a



COBE measured: $v = 371 \pm 1 \text{ km s}^{-1}$, which gives T differences of about 3 mK in each direction.

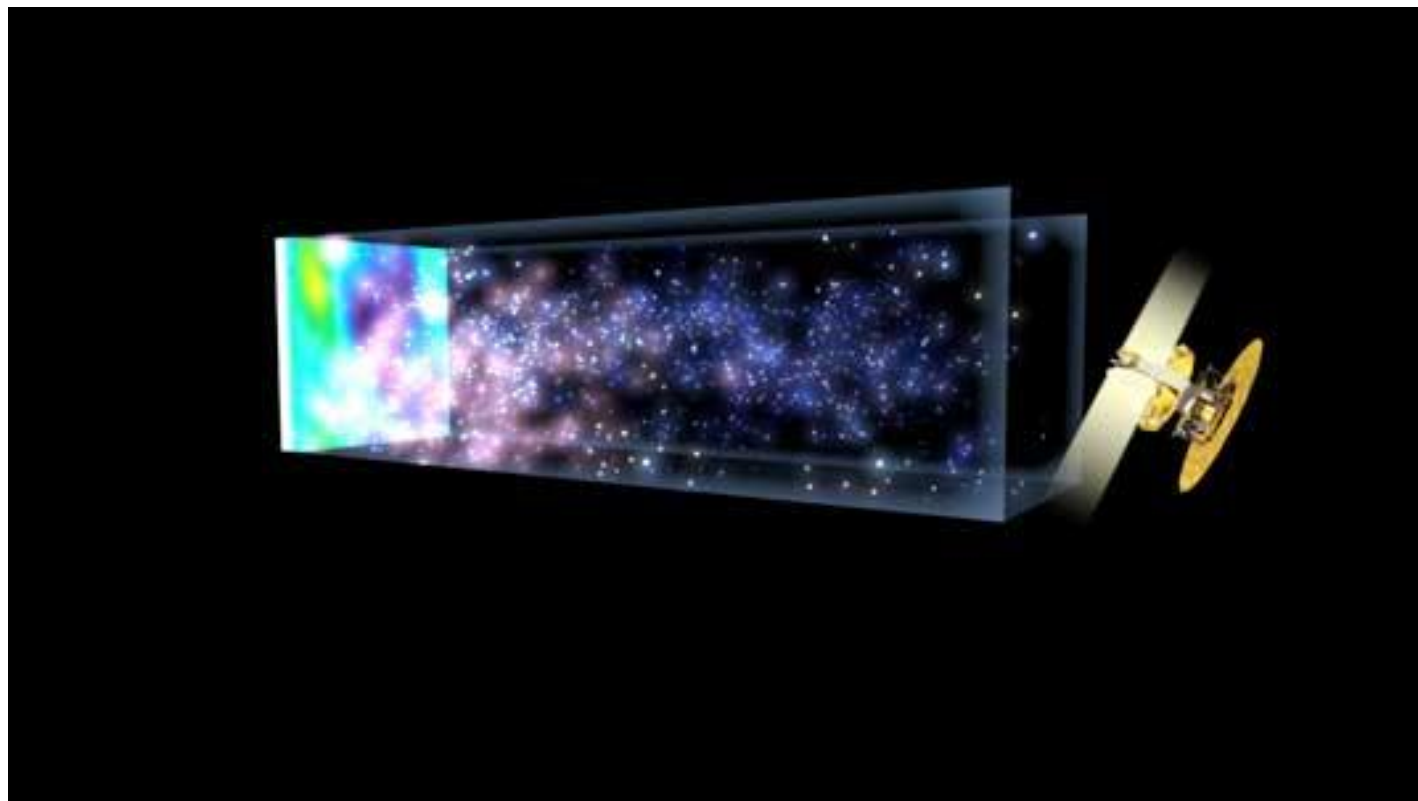
ANIZOTROPIJA 1 u 100 000

The view from COBE

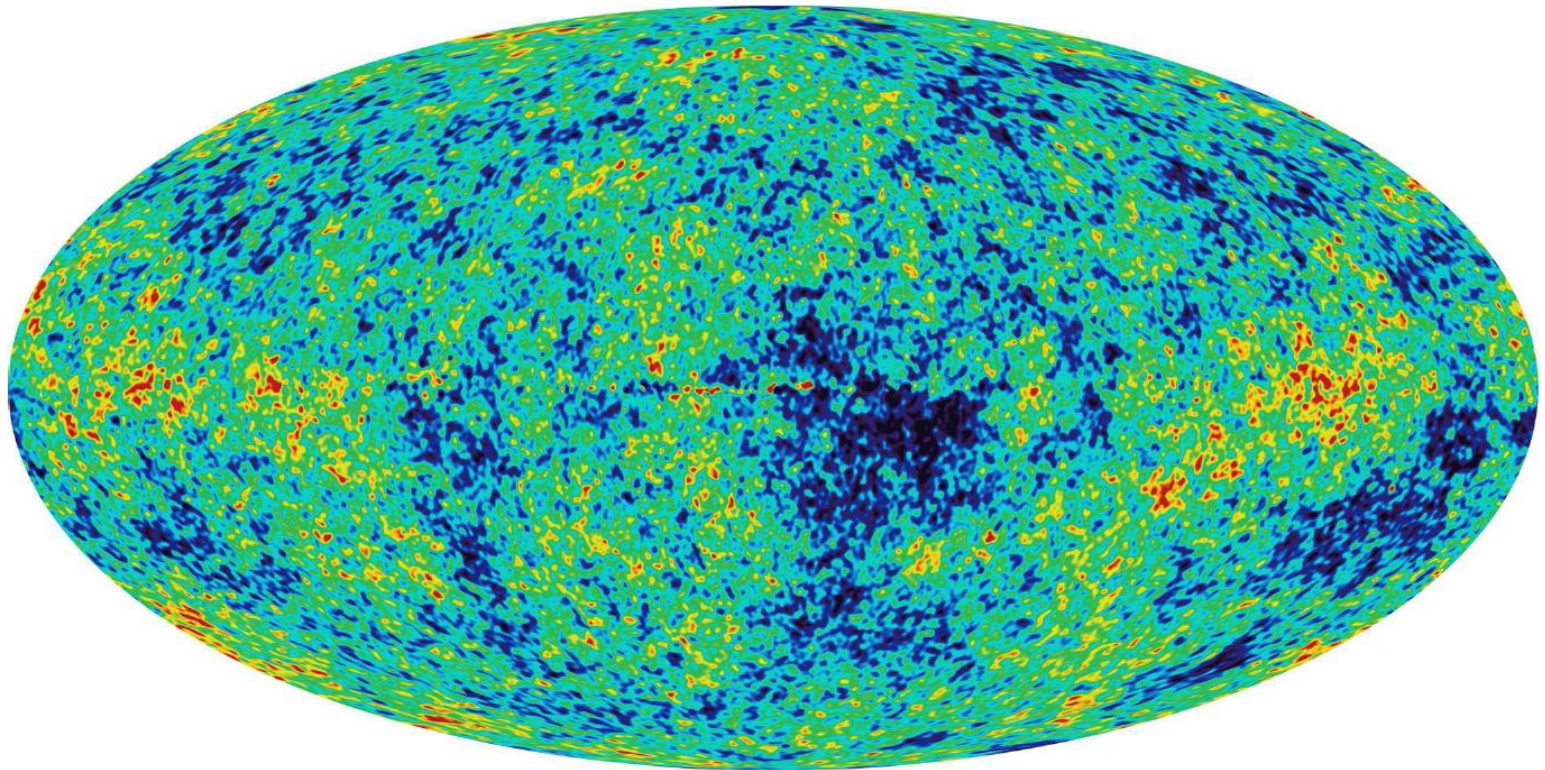


The map is in galactic co-ordinates. The red stripe is emission from the galactic plane.

PORIJEKLO STRUKTURA NA NAJVEĆOJ SKALI – u kvantnim fluktuacijama sa samog početka svemira



ANIZOTROPIJA nastala na $z=1100$ -sjeme formiranja opaženih struktura u svemiru



Spektar intenziteta zračenja – (“power spectrum”), gdje se temp. fluktuacija oko srednje temp. u smjeru α izražava autokorelacijskom funkcijom:

$$C(\theta) = \left\langle \frac{\delta T}{T}(\alpha) \frac{\delta T}{T}(\alpha + \theta) \right\rangle$$

MULTIPOLNI RAZVOJ

$$C(\theta) = \frac{1}{4\pi} \sum_{l=2}^{\infty} a_l^2 (2l + 1) P_l(\cos \theta)$$

- $l=0$ – srednja temperatura neba
- $l=1$ – dipolna anizotropija
- Viši multipoli odgovaraju fluktuacijama kutnih veličina:

$$\theta \approx \frac{180^\circ}{l}$$

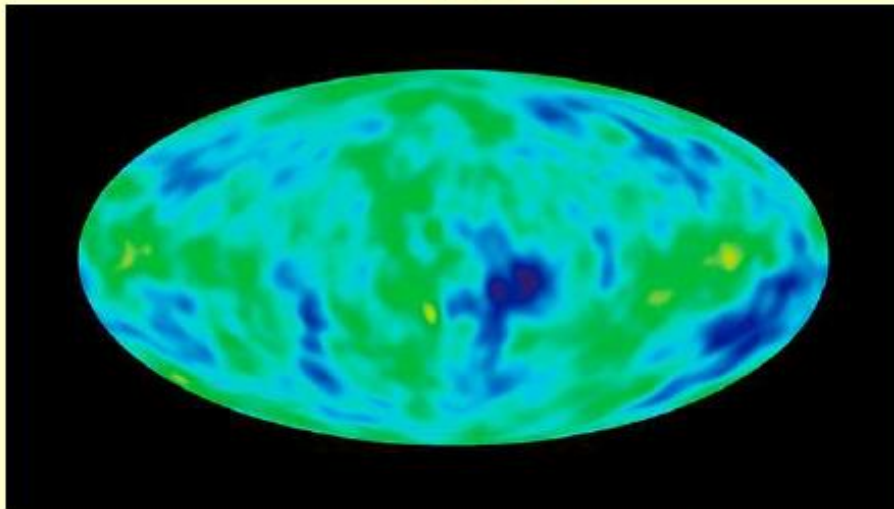
Spektar intenziteta fluktuacija

The powers a_l^2 are adjusted to give a best fit of $C(\theta)$ to the observed temperature distribution. The resulting distribution of a_l^2 values versus l is called the **Power Spectrum** of the fluctuations.

Higher angular resolution means that more terms of high l must be included.

ANIZOTROPIJE CMB-a

COBE (1993)

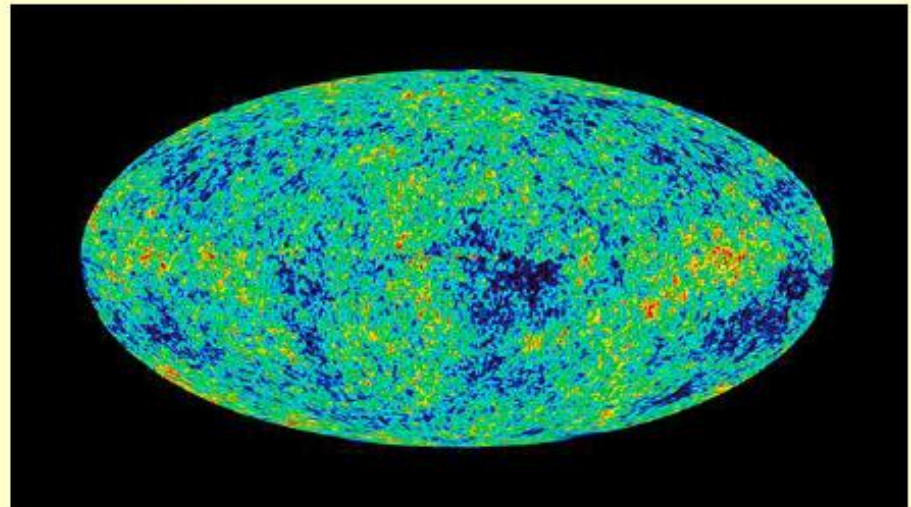


*Sky map, all
foregrounds
subtracted*

WMAP 7 years (2008)

- CMB: anisotropies are very small:

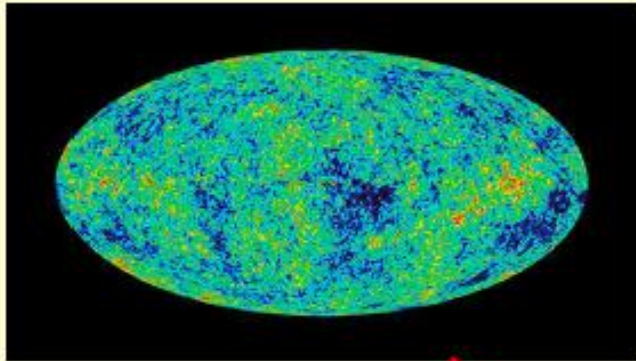
$$\Delta T \sim 18 \mu K \Rightarrow \Delta T / T \sim 10^{-5}$$



@WMAP web site

Analysis of T anisotropies

- Angular fluctuation spectrum: strength of T fluctuations (hot and cold spots) as a function of angular size

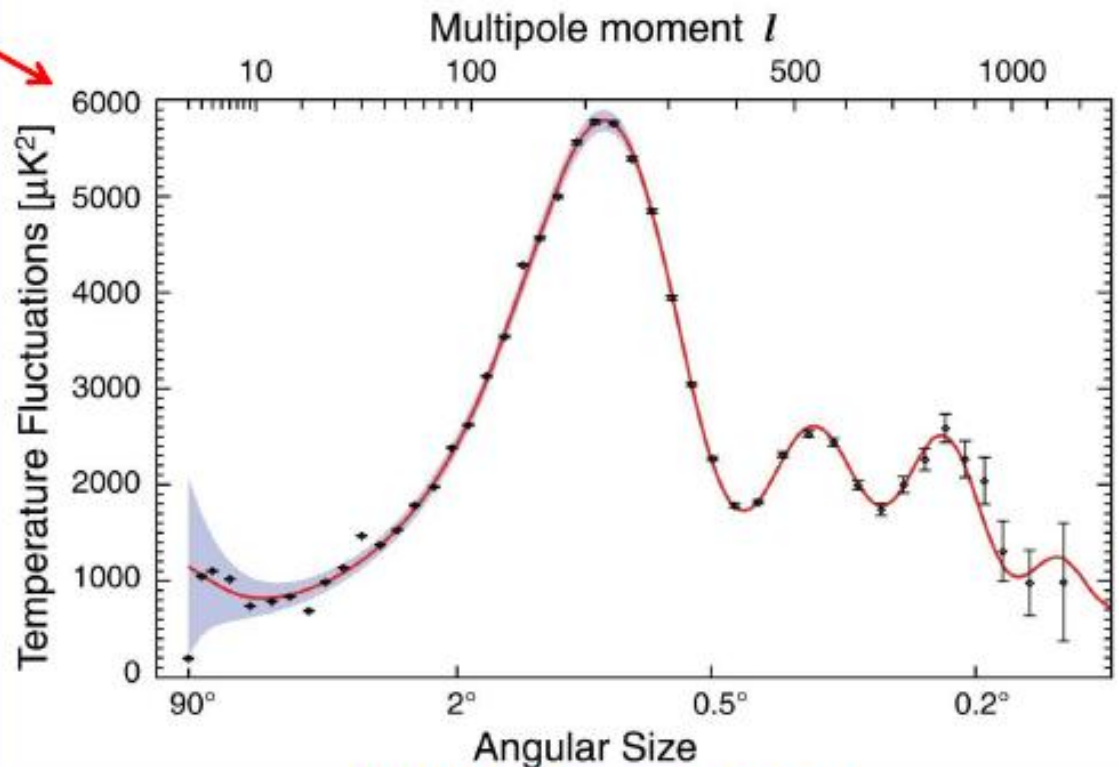


main T fluctuations:
angular size of $\sim 1^\circ$
($l \sim 200$)

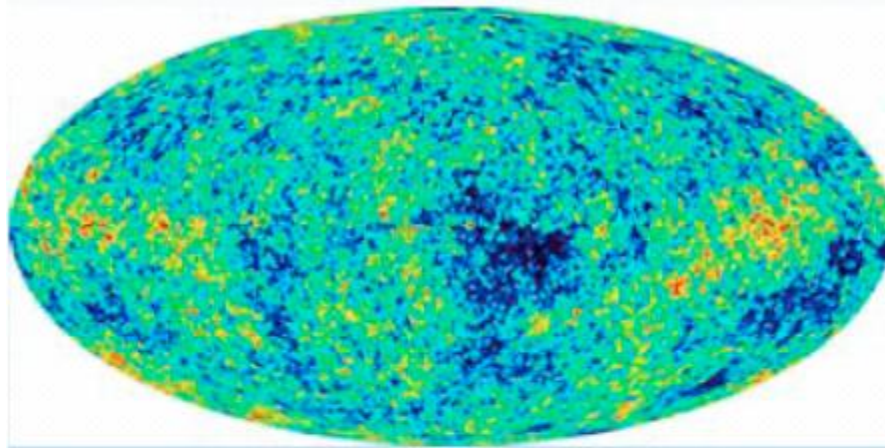
$$\Delta\theta \approx \frac{\pi}{l} \approx \frac{s_*}{D_*}$$

s_* = distance traveled by sound since $t=0$ (sound horizon)

D_* = comoving distance to z_*
(angular diameter distance)



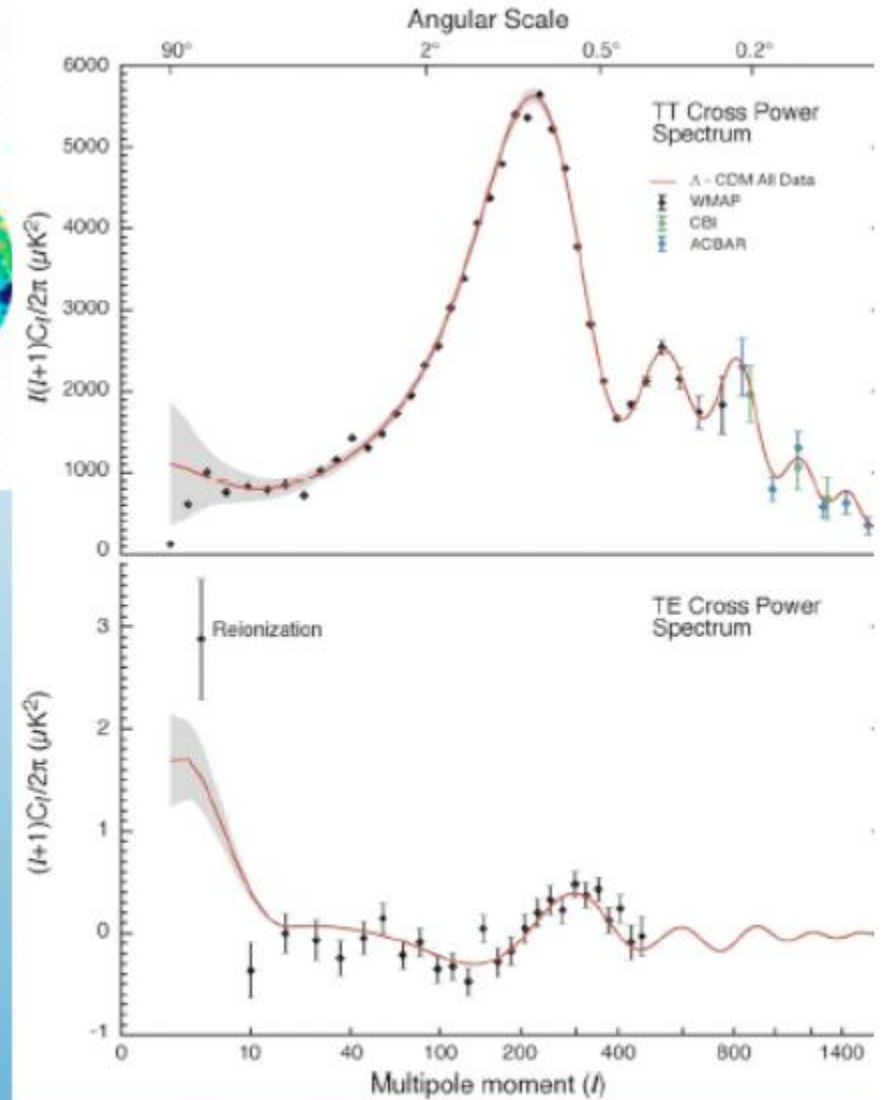
WMAP 7 years (2008)



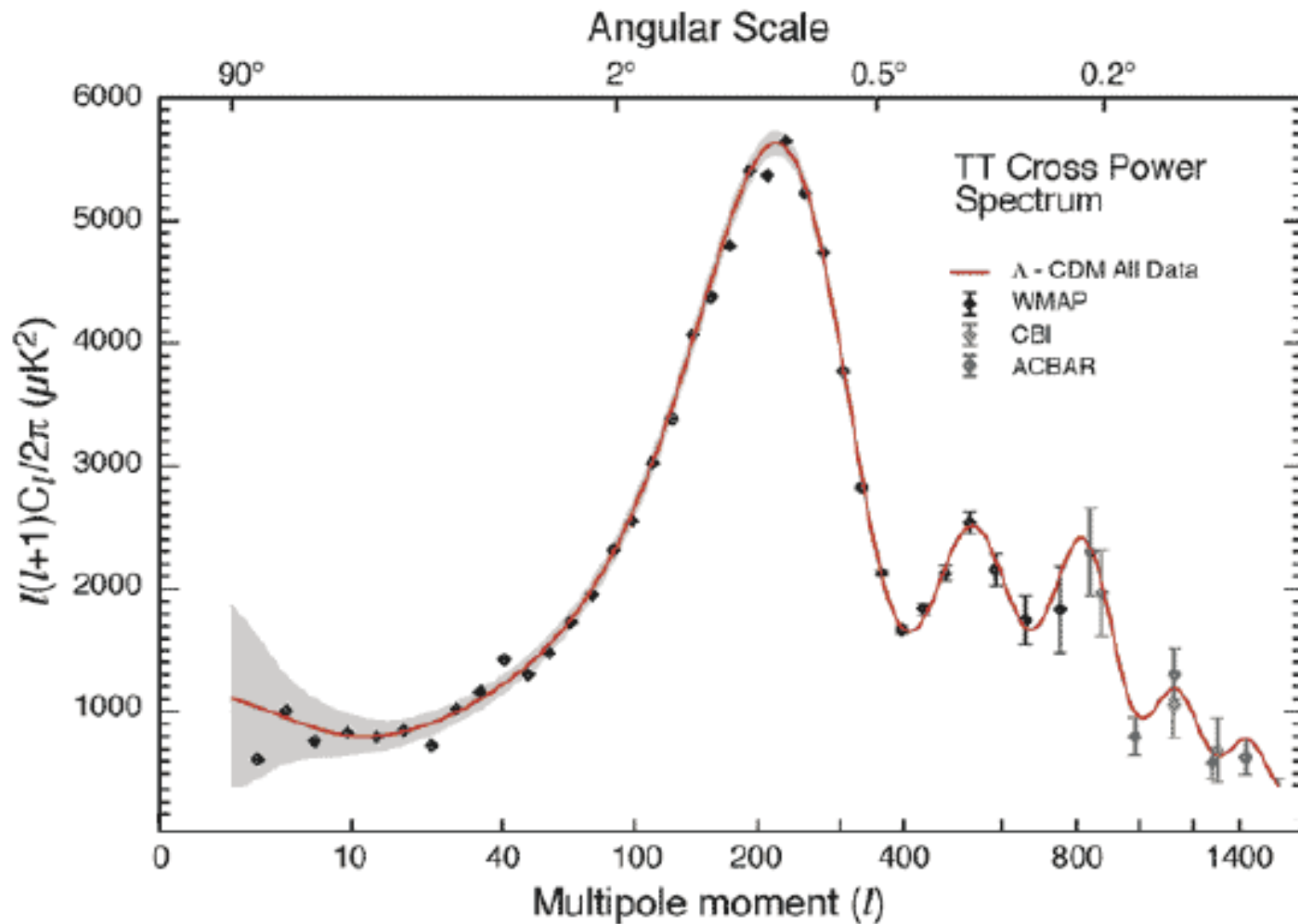
WMAP best fit
 (WMAPext + 2dFGRS +
 Lyman α +running sp.
 index)

$$\Omega_B h^2 = 0.0224 \pm 0.0009$$

$$\eta_{10} = 6.14 \pm 0.25$$

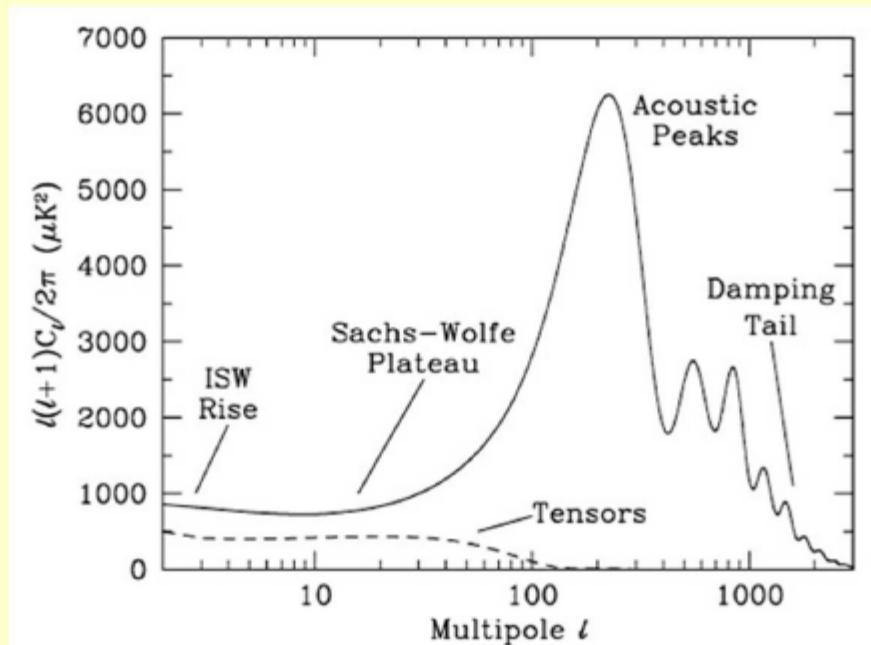


Rezultati WMAPA-a

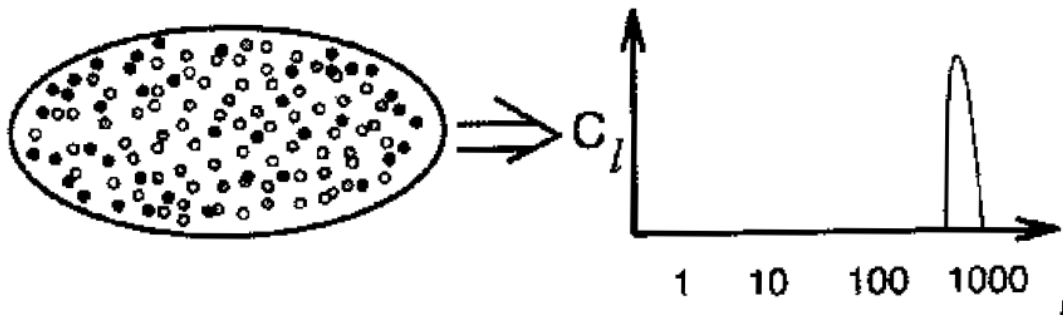
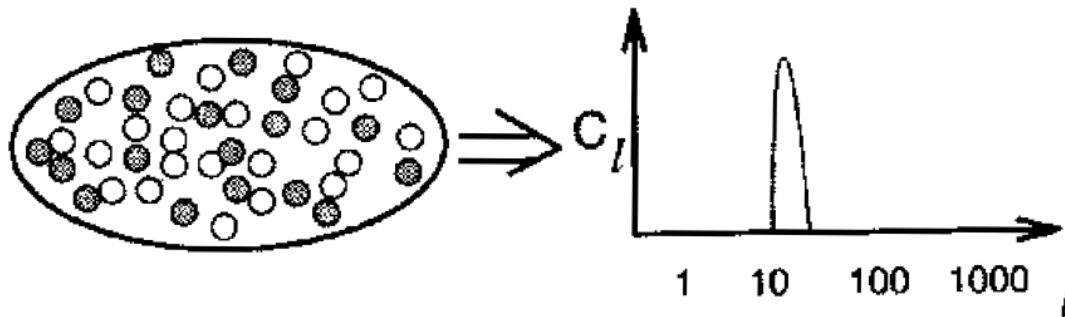
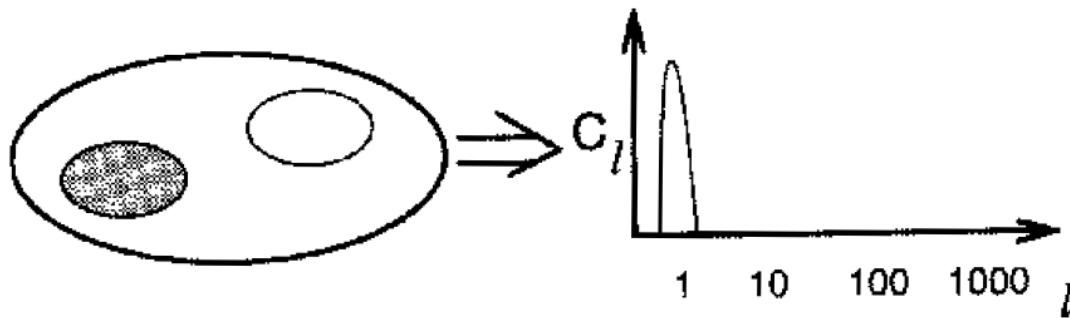


PORIJEKLO ANIZOTROPIJA

- variations in the density or velocity of the plasma at the last scattering surface: acoustic peaks
- variations in the gravitational potential at the LSS or along the photon path: low multipole plateau (large scales)



KARAKTERISTIČNE SKALE ANIZOTROPIJA CMB-a



IZVORI ANIZOTROPIJA CMB-a

koje nastaju u vrijeme rekombinacije

- **GRAVITACIJSKE/Sachs-Wolfeove** perturbacije
- **INTRISIČNE/adijabatske** perturbacije
- **DOPPLEROVE** perturbacije

Sachs-Wolfeov efekt

- **GRAVITACIJSKE/Sachs-Wolfeove** perturbacije: fotoni koji izlaze iz područja veće (**manje**) gustoće, doživljavaju crveni (**plavi**) pomak pri penjanju uz (**kotrljanju niz**) gravitacijski potencijal. Gravitacijski crveni pomak (praćen dilatacijom vremena) vodi na fluktuacije (do 3 stupnja) $\delta T/T$ koje su linearne u fluktuacijama gustoće $\delta\rho/\rho$:

$$\frac{\delta T}{T} \approx -\frac{1}{3} \left(\frac{L}{ct} \right)^2 \frac{\delta\rho}{\rho}$$

Intrisične perturbacije

- U područjima velike gustoće zračenje je komprimirano i time dodatno ugrijano. Suptilnost je da gušća područja rekombiniraju kasnije – imaju manji crveni pomak i pojavljuju se kao vruća. Opažene temp. perturbacije:

$$\left(\frac{\delta T}{T}\right)_{\text{obs}} = \frac{\delta \rho}{\rho}$$

Dopplerove perturbacije

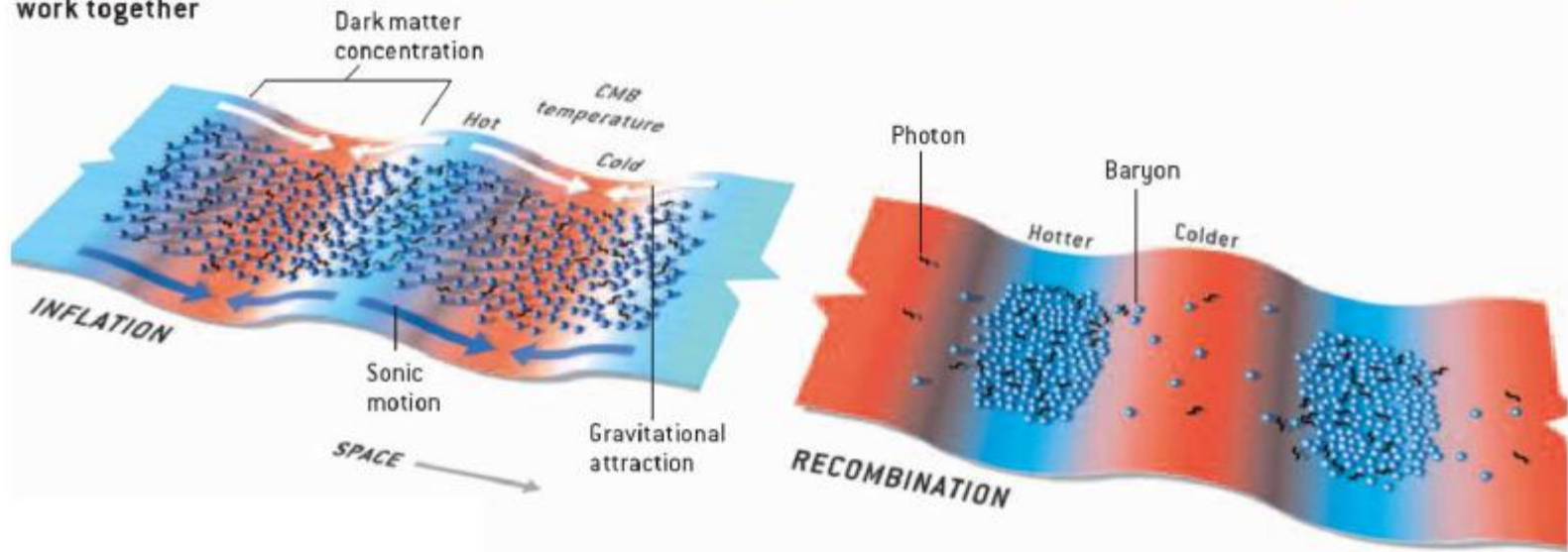
- **Konačna brzina plazme** u vrijeme rekombinacije daje Dopplerove pomake u frekvenciji, a time i u temperaturi. Taj se efekt opaža na najmanjoj kutnoj skali

PORIJEKLO ANIZOTROPIJA

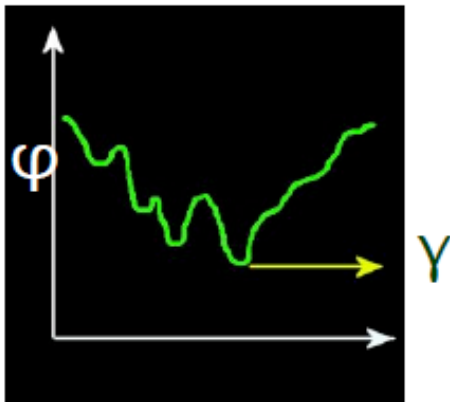
- **dark matter** collapse creates **over-densities**, making the photon-baryon plasma to **oscillate** into and out of these over-densities.
- At recombination: oscillation pattern is **frozen** \Rightarrow imprint in CMB spectrum (and matter spectrum)

FIRST PEAK

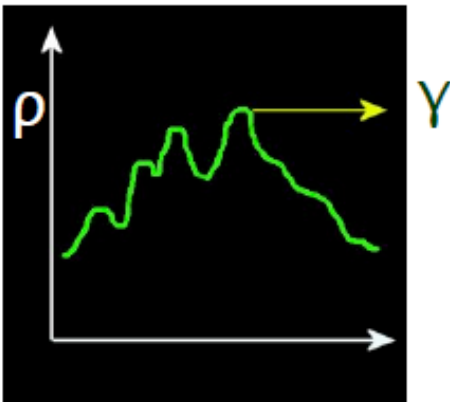
Gravity and sonic motion work together



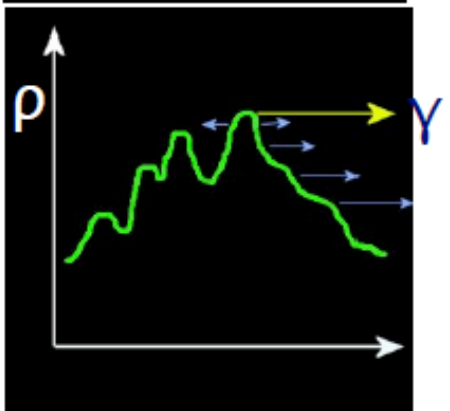
Fluctuations in the matter density \rightarrow fluctuations in the CMB temperature



Photons are **redshifted** as they move out of gravitational potential wells



Dense regions have higher temperature
 \Rightarrow photons have higher energy



Photons emitted from a moving surface are **red/blue-shifted**

Fortunately the effects do not *quite* cancel so the CMB carries a memory of the past

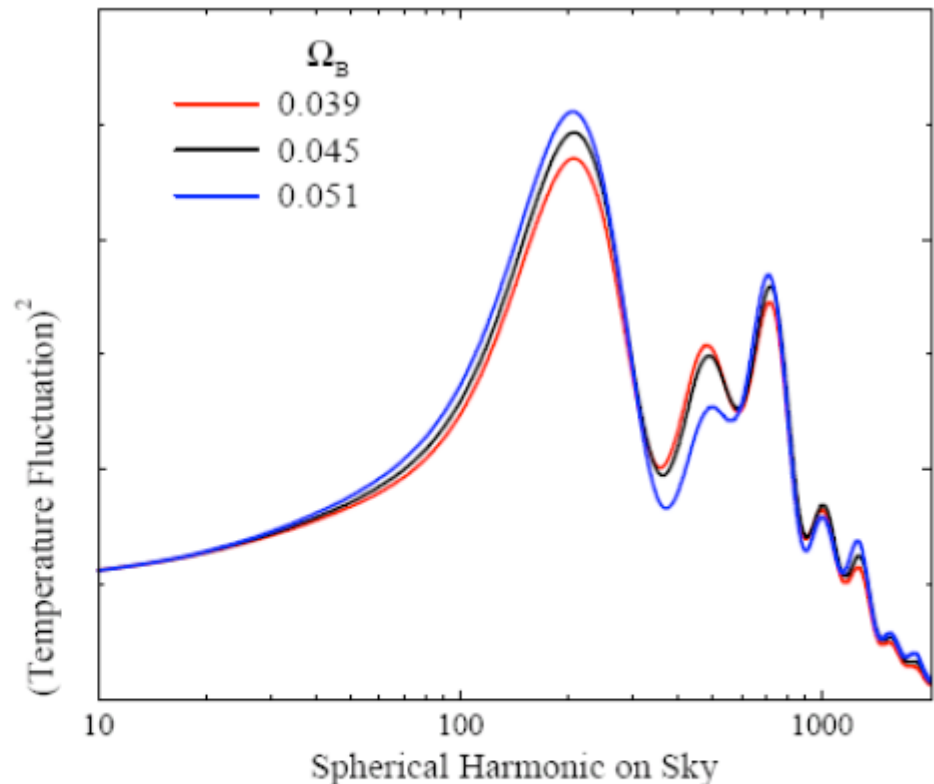
CMB KAO "BARYOMETAR"

ΔT_ℓ provide *independent* measure of $\Omega_B h^2$

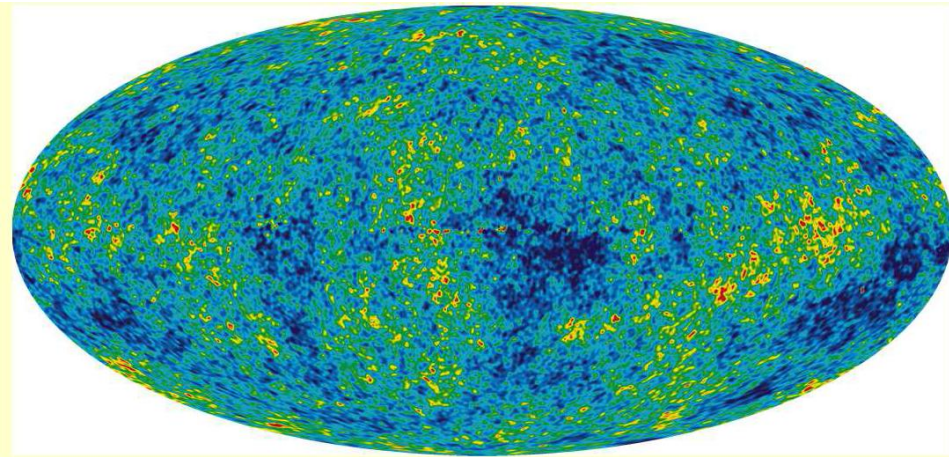
Acoustic oscillations in (coupled photon-baryon fluids) Imprint features at $< 1^\circ$ in angular power spectrum

Peak positions and heights sensitive to cosmological parameters e.g.

Ratio of 2nd peak/1st peak
 \Rightarrow **baryon density**



BBN vs CMB determinations of baryon density \rightarrow fundamental test of cosmology and thermal history at $z \sim 10^3 - 10^{10}$



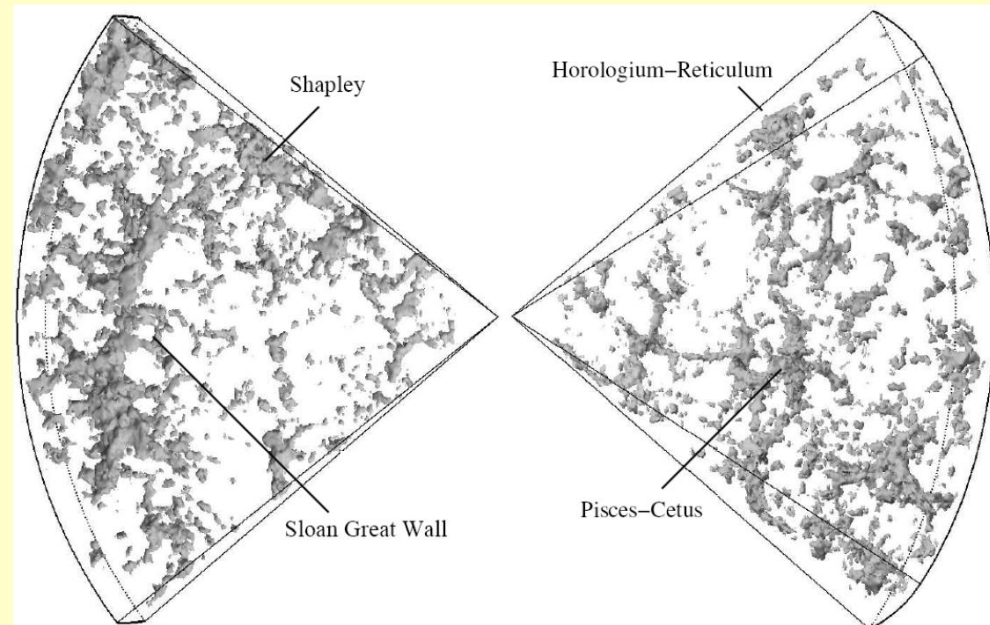
WMAP (2008), all foregrounds subtracted

- CMB: very small anisotropies

$$\Delta T/T = 10^{-5}$$

- matter density inhomogeneities, amplified by gravitation after decoupling → structures

- Large scale structures today: galaxies, clusters, superclusters, voids and filaments
- Formation of structures: cold dark matter mandatory



2dF Galaxy Redshift Survey, (2007) 21

A possible scenario

